

SPHERICALLY SYMMETRIC
VACUUM SOLUTIONS IN
FIRST ORDER GRAVITY:
FATE OF
CURVATURE SINGULARITIES

SANDIPAN SENGUPTA

IIT KHARAGPUR

INDIA

①

● GENERAL RELATIVITY IS FAMOUSLY KNOWN TO EXHIBIT BLACK HOLE SOLUTIONS

● SUCH CONFIGURATIONS ARE ASSOCIATED WITH CURVATURE SINGULARITIES

● THE SIMPLEST EXAMPLE OF A BH SOLUTION IS PROVIDED BY THE SCHWARZSCHILD METRIC: (SPHERICALLY SYMMETRIC, STATIC) :

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{2M}{r}\right)} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

● THIS DESCRIBES THE SPACETIME GEOMETRY OUTSIDE A (SPH. SYM) MASS M ; CHARACTERISED BY A SURFACE AT $r = 2M$ WHICH HAS A DEGENERATE 3-METRIC AND DIVIDES SPACETIME INTO AN INTERIOR AND AN EXTERIOR REGION ($r < 2M$) ($r > 2M$)

- THE SCHW. METRIC CONTINUES TO BE A SOLUTION OF EINSTEIN EQ^N (IN VACUUM) $R_{\mu\nu} = 0$ EVEN AT THE INTERIOR

- IF ONE CONTINUES INSIDE, A CURVATURE SINGULARITY AT $r=0$ IS ENCOUNTERED, AS THE $SO(3,1)$ FIELD STRENGTH COMPONENTS $R_{\mu\nu}{}^{15}(\omega)$ AS WELL AS THE SCALAR $R_{\mu\nu}{}^{15} R^{\mu\nu}{}_{15} = R_{\mu\nu\alpha\beta} R^{\mu\nu\alpha\beta} = \frac{48M^2}{r^6}$ BLOW UP

- WHEN THE METRIC IS INVERTIBLE ($\text{Det} g_{\mu\nu} \neq 0$), THERE SEEMS TO BE NO ESCAPE FROM SUCH A SCENARIO IN GR

- POSSIBLE WAY OUT TO EVADE SINGULARITIES: MODIFIED GRAVITY, EXTRA DIMENSIONS, EXOTIC MATTER, QUANTUM EFFECTS, ...

3

- HOWEVER, ONE CAN STILL CHOOSE TO REMAIN CONSERVATIVE:
IS IT POSSIBLE TO CONSTRUCT A WELL-BEHAVED INTERIOR GEOMETRY IN PURE GRAVITY WITHOUT MODIFYING THE STANDARD THEORY (ACTION PRINCIPLE / LAGRANGIAN)?
- IN THAT CASE, THE ONLY POSSIBLE OPTION (PERHAPS) IS TO WORK WITH DEGENERATE (NON-INVERTIBLE) METRIC AT THE INTERIOR!

VACUUM SOLUTIONS WITH DEGENERATE INTERIOR METRIC

- SUCH A SOLUTION, IF IT EXISTS, SHOULD ALSO EXHIBIT THE FOLLOWING PROPERTIES:
 1. BOTH THE (NON-DEGENERATE) EXTERIOR METRIC AND THE (DEGENERATE) INTERIOR METRIC MUST BE SOLUTIONS OF VACUUM EOM IN GRAVITY THEORY
 2. THE INTERIOR & EXTERIOR GEOMETRIES MUST BE JOINED SUFFICIENTLY SMOOTHLY AT THE 'HORIZON'
 3. THE CONFIGURATION MUST BE 'REGULAR' EVERYWHERE

(4)

CONSTRUCTION OF EXTERIOR SOLUTION :

- LET US CONSIDER A SPH. SYM. STATIC METRIC OF THE FOLLOWING FORM (FOR THE EXTERIOR REGION) :

$$ds^2 = - \left(1 - \frac{2M}{f(u)}\right) dt^2 + \frac{f'^2(u)}{\left(1 - \frac{2M}{f(u)}\right)} du^2 + f^2(u) d\Omega_{\theta, \phi}^2$$

$$[-\infty < u < \infty]$$

- $f(u) = 2M \left[1 - \left(1 - \frac{u}{u_0}\right)^{2l+1}\right]$, $l \gg 1$

$$f(u_0) = 2M, \quad f'(u_0) = 0$$

$$\Rightarrow \underline{\text{Det } g_{\mu\nu}(u_0) = 0}$$

- THIS METRIC IS GAUGE-EQUIVALENT TO THE SCHW. METRIC AT $u > u_0$, BUT NOT AT $u = u_0$

- $u < u_0$, $u > u_0$, $u = u_0$ & $u = 0$ CORRESPOND TO $r < 2M$, $r > 2M$, $r = 2M$ & $r = 0$, RESPECTIVELY; HOWEVER, THE REGION $-\infty < u < 0$ HAS NO ANALOGUE IN SCHW. COORDINATES.

(5)

- AT THE 'HORIZON' $u = u_0$, THE $SO(3,1)$ FIELD-STRENGTH COMPONENTS

$$R_{MN}{}^{IJ} = \partial_{[M} W_{N]}{}^{IJ} + W_{[M}{}^{IK} W_{N]K}{}^J \quad \text{BECOME:}$$

$$R^{01} \stackrel{*}{=} R^{02} \stackrel{*}{=} R^{03} \stackrel{*}{=} 0,$$

$$R^{12} \stackrel{*}{=} R^{31} \stackrel{*}{=} 0, \quad R^{23} \stackrel{*}{=} \sin\theta \, d\theta \wedge d\phi$$

- TO EMPHASIZE, ~~THE~~ THE COMPONENTS OF g_{MN} AND $R_{MN}{}^{IJ}$ AT $u = u_0$ FOR THE EXTERIOR GEOMETRY MUST MATCH THOSE FOR THE INTERIOR GEOMETRY, WHICH WE CONSTRUCT NEXT.

⑥ CONSTRUCTION OF INTERIOR: DEGENERATE METRIC

- THE AIM IS TO CONSTRUCT AN INTERIOR ($u < u_0$) GEOMETRY THAT CONNECTS SMOOTHLY TO THE (DEGENERATE) HORIZON, AND IS BASED ON A NON-INVERTIBLE METRIC.
- THIS MUST BE A SOLUTION OF THE EOMS IN PURE GRAVITY
- THESE EOMS MUST COME FROM A LAGRANGIAN THAT CAN ACCOMODATE NON-INVERTIBLE METRICS
- HENCE, THE APPROPRIATE LAGRANGIAN (DENSITY) MUST CORRESPOND TO THE FIRST ORDER FORMULATION, WHICH DOES NOT REQUIRE THE INVERSE METRIC IN ITS CONSTRUCTION:

$$\mathcal{L}(e, \omega) = \frac{1}{k^2} \epsilon^{\mu\nu\alpha\beta} \epsilon_{\gamma\delta\kappa\lambda} e_{\mu}^{\gamma} e_{\nu}^{\delta} R_{\alpha\beta}^{\kappa\lambda}(\omega)$$

$$S = \frac{1}{k^2} \int d^4x \mathcal{L}(e, \omega)$$

[HILBERT-PALATINI]

(7)

DEGENERATE SOLUTIONS: NON-EINSTEINIAN GRAVITY

- THIS IS TO BE CONTRASTED WITH THE SECOND ORDER LAGRANGIAN (DENSITY), WHICH IS ILL-DEFINED FOR DEGENERATE METRICS, AND DOES NOT ACCOMMODATE DEGENERATE SOLUTIONS:

$$\mathcal{L}(g) = \frac{1}{k^2} \sqrt{-g} g^{\mu\nu} R_{\mu\nu}(g)$$

- THE EOMS, COMING FROM THE VARIATIONS OF $\mathcal{L}(e, \omega)$ W.R.T. e_M^I & ω_M^{IJ} , ARE:

$$\delta \omega: e_{[M}^I \nabla_\nu e_{\alpha J]} = 0$$

$$\delta e: e_{[M}^I R_{\alpha\beta J]} = 0$$

- WHEN $\text{Det } e_M^I \neq 0$, THESE EOM-S CORRESPOND TO EINSTEIN'S THEORY

- WHEN $\text{Det } e_M^I = 0$, ONE DOES NOT RECOVER EINSTEIN'S THEORY

(8)

- SOLUTIONS OF THE FIRST-ORDER EOM WITH DEGENERATE METRICS GENERICALLY EXHIBIT TORSION; THESE ARE NOT PERCEIVED AT ALL BY THE SECOND-ORDER THEORY BASED ON INVERTIBLE $g_{\mu\nu}$.

[TSEYTLIN 1982, KAUL & SENGUPTA 2016]

- HOW TO CONSTRUCT A DEGENERATE SPACETIME SOLUTION (INTERIOR) IN THE PRESENT CONTEXT OF A SPHERICALLY SYMMETRIC GEOMETRY?

(IN ADDITION,)
CAN ONE OBTAIN WELL-BEHAVED $SO(3,1)$ FIELD-STRENGTH COMPONENTS EVERYWHERE, AND ALSO SATISFY THE REQUIREMENT OF CONTINUITY PROPERTIES AT THE 'HORIZON' $u = u_0$?

(9)

INTERIOR: AN EXPLICIT DEGENERATE SOLUTION

- FOR $u \leq u_0$, CONSIDER THE DEGENERATE METRIC ($g_{tt}=0$) WITH TOPOLOGY $\mathbb{R} \times S^2$:

$$ds_{(4)}^2 = 0 + F^2(u) du^2 + H^2(u) [d\theta^2 + \sin^2\theta d\phi^2]$$

- TETRAD:

$$e_{\mu}^I = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & F(u) & 0 & 0 \\ 0 & 0 & H(u) & 0 \\ 0 & 0 & 0 & H(u) \sin\theta \end{pmatrix} \equiv \begin{pmatrix} 0 & 0 \\ 0 & e_a^i \end{pmatrix}$$

- SPIN-CONNECTION:

$$\omega^{0i} = 0, \quad \omega^{ij} = \bar{\omega}^{ij}(e) + K^{ij}$$

$$[I \equiv (0, i) \equiv (0, 1, 2, 3)]$$

- MOST GENERAL K^{ij} ~~SATISFYING~~ SATISFYING THE EOMS: $K^{ij} = \epsilon^{ijk} e_a^1 N_{ke}$ Where

$$N_{ke} = N_{ek} \equiv \begin{pmatrix} \alpha & \eta_3 & \eta_2 \\ \eta_3 & \beta & \eta_1 \\ \eta_2 & \eta_1 & \gamma \end{pmatrix}$$

(10)

- THE FIELDS $e_\mu^I \equiv (F(u), H(u))$ and $\omega_\mu^{15} \equiv (\bar{\omega}_\mu^{15}(e), K_\mu^{15})$ SATISFY THE 1ST ORDER EOM-s, PROVIDED THEY ARE RELATED THROUGH THE CONSTRAINT:

$$\eta_1^2 + \eta_2^2 + \eta_3^2 - \alpha\beta - \beta\gamma - \gamma\alpha = \frac{H^{12}}{F^2 H^2} + \frac{2}{FH} \left(\frac{H'}{F}\right)' - \frac{1}{H^2}$$

$$\Rightarrow C(\eta_1, \eta_2, \eta_3, \alpha, \beta, \gamma, F, H) = 0$$

- FOR SIMPLICITY, WE CONSIDER 'MINIMAL' TORSION:

ONLY $\eta_1 \neq 0$. AND ASSUME THAT $\eta_1 = \eta_1(u)$ [SPH. SYM].

- THE ABOVE CONSTRAINT BECOMES:

$$C(\eta_1, F, H) = 0$$

- NEED 2 FURTHER CONDITIONS TO DETERMINE η_1, F & H .

(11)

INTERIOR GEOMETRY :

- THE $SO(3,1)$ FIELD-STRENGTH COMPONENTS:

$$R^{12} = \left(\eta_1 H - \frac{H'}{F} \right)' du \wedge d\theta$$

$$R^{23} = \left[1 + \left(\eta_1 H - \frac{H'}{F} \right) \left(\eta_1 H + \frac{H'}{F} \right) \right] \sin\theta d\theta \wedge d\phi$$

$$R^{31} = \left(\eta_1 H + \frac{H'}{F} \right)' \sin\theta du \wedge d\phi + 2\eta_1 H \cos\theta d\theta \wedge d\phi$$

$$R^{0i} = 0$$

- CONTINUITY OF $R_{\mu\nu}$ AT $u = u_0$ (HORIZON)

IMPLIES:

$$R^{12}(u_0) = 0, \quad R^{23}(u_0) = \sin\theta d\theta \wedge d\phi, \quad R^{31}(u_0) = 0,$$

$$R^{0i}(u_0) = 0$$

- THE 2 CONDITIONS TO BE CHOSEN MUST BE CONSISTENT WITH THE ABOVE; WE CHOOSE:

$$\eta_1 H = \frac{H'}{F},$$

$$\frac{H'}{F} = \beta e^{\alpha(u-u_0)} (u_0 - u)^m, \quad m > 1$$

(12)

• THESE IMPLY: $R_{\mu\nu}^{(INT)} = R_{\mu\nu}^{(EXT)}$ AT $u = u_0$

• THE 3 UNKNOWN η_1, F, H CAN NOW BE SOLVED USING THE 3 EQUATIONS:

$$\eta_1 = \eta_1(u), \quad F = F(u), \quad H = H(u).$$

•
$$ds_{INT}^2 = 0 + F^2(u) du^2 + H^2(u) (d\theta^2 + \sin^2\theta d\phi^2)$$

$$F(u) \rightarrow 0 \quad \text{as } u \rightarrow u_0 \quad \text{or } u \rightarrow -\infty$$

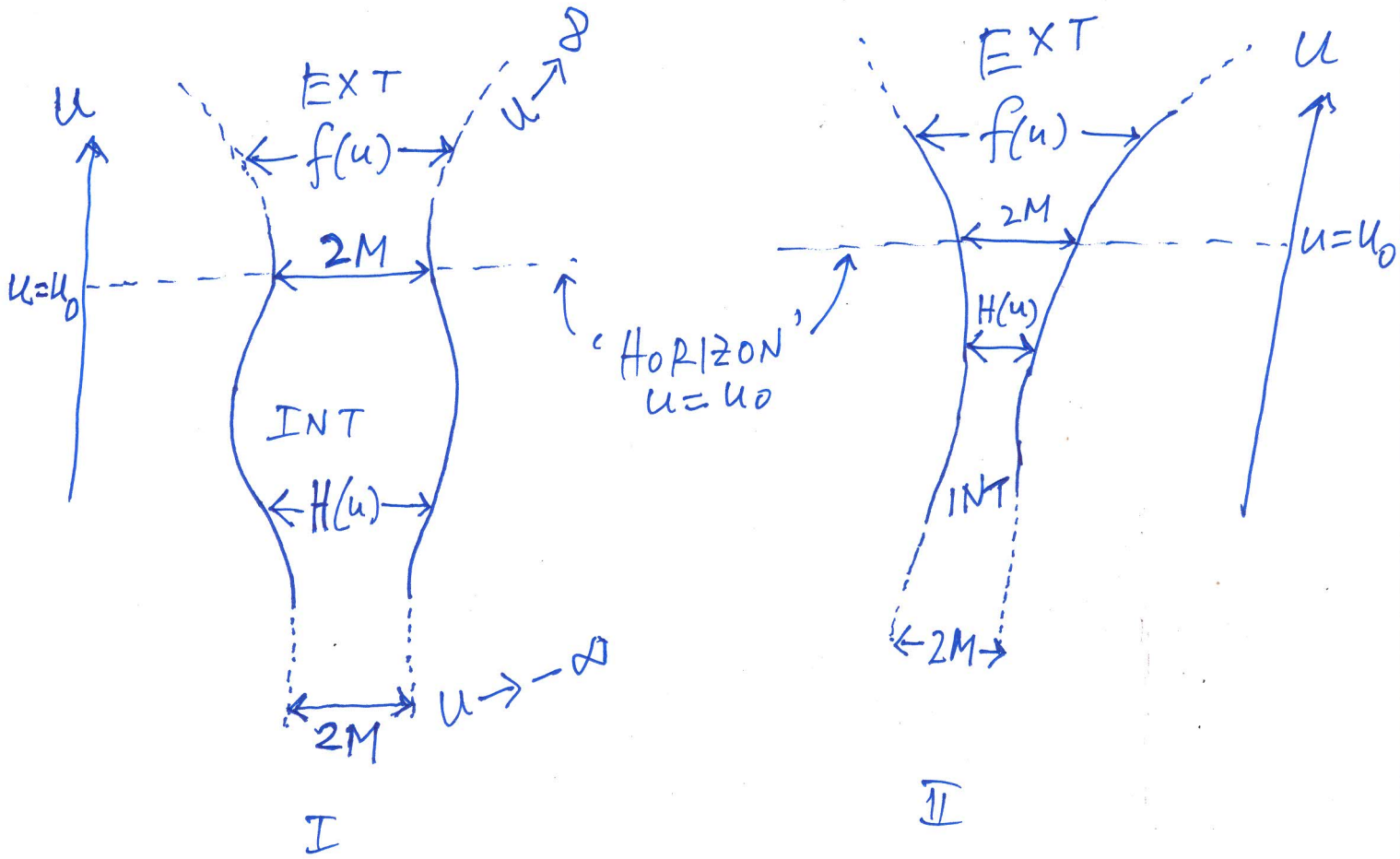
$$H(u) \rightarrow 2M \quad \text{as } u \rightarrow u_0 \quad \text{or } u \rightarrow -\infty$$

$$\eta_1(u) \rightarrow 0 \quad \text{as } u \rightarrow u_0 \quad \text{or } u \rightarrow -\infty$$

• THIS, IN TURN, IMPLIES THAT:
$$g_{\mu\nu}^{(INT)} = g_{\mu\nu}^{(EXT)} \quad \text{AT } u = u_0$$

• ALL THE $SO(3,1)$ FIELD-STRENGTH COMPONENTS $R_{\mu\nu}$ ARE FINITE AT $-\infty < u < \infty$ [UNLIKE SCHW. INTERIOR]

PICTORIAL REPRESENTATION :
(INTERIOR + EXTERIOR \equiv FULL SOLUTION)



TO SUMMARISE

- A NEW CLASS OF VACUUM SOLUTIONS (OF FIRST ORDER EOM-s) HAS BEEN CONSTRUCTED ; HAS INVERTIBLE METRIC AT THE EXTERIOR & NON-INVERTIBLE METRIC AT THE HORIZON AND THE INTERIOR
- THE SOLUTION IS 'REGULAR' EVERYWHERE , SINCE THE $SO(3,1)$ FIELD-STRENGTH COMPONENTS ARE FINITE EVERYWHERE
- THIS IS UNLIKE THE ~~CASE~~ CASE OF INTERIOR SCHW. SOLUTION DESCRIBING A BLACK HOLE , WHERE THESE ARE ILL-DEFINED AT THE ORIGIN $r=0$.

(15)

- THE METRIC & FIELD-STRENGTH COMPONENTS ARE CONTINUOUS AT THE HORIZON
 - THE DEGENERATE INTERIOR SPACETIME SOURCES TORSION ~~IS~~
 - THE ORIGIN OF TORSION IN THIS CONFIGURATION IS PURELY GEOMETRIC (NOT MATTER-COUPLING)
 - THE CLASS OF SOLUTIONS CORRESPOND TO 'VANISHING (FINITE!) ACTION; SHOULD BE TREATED AS SADDLE POINTS IN PATH INTEGRAL
- QUANTUM GRAVITY

(16)

TO EXPLORE FURTHER

- ARE THESE SOLUTIONS REALLY NON-SINGULAR?



BEHAVIOUR OF GEODESIC CONGRUENCES

- THE TORSIONAL FIELD CONTENT IS CONFINED AT THE INTERIOR;

HOW TO MAKE USE OF THIS ADDITIONAL FIELD CONTENT?



MANY EXCITING POSSIBILITIES

- HOW TO MODEL STARS (INTERIOR WITH MATTER.) OR OTHER POSSIBLE REALISTIC SYSTEMS?

THANK

YOU!